

Theory vs molecular modelling of charge transfer reactions: some problems and challenges



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Outline

1. Motivation
2. Medium (solvent) coordinates
3. Quantum effect of solvent on activation barrier
4. Orbital overlap (electronic transmission coefficient)
5. Some problems ahead

Transition state theory (TST)

saddle point

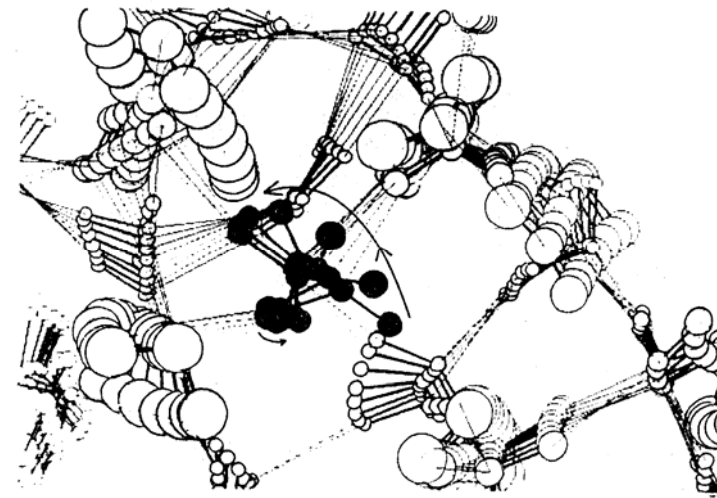


**Three-dimensional reaction energy surface
(solvent and intramolecular coordinates)**

A simple way to define the solvent coordinate

(Marcus theory)

λ is the solvent reorganization energy



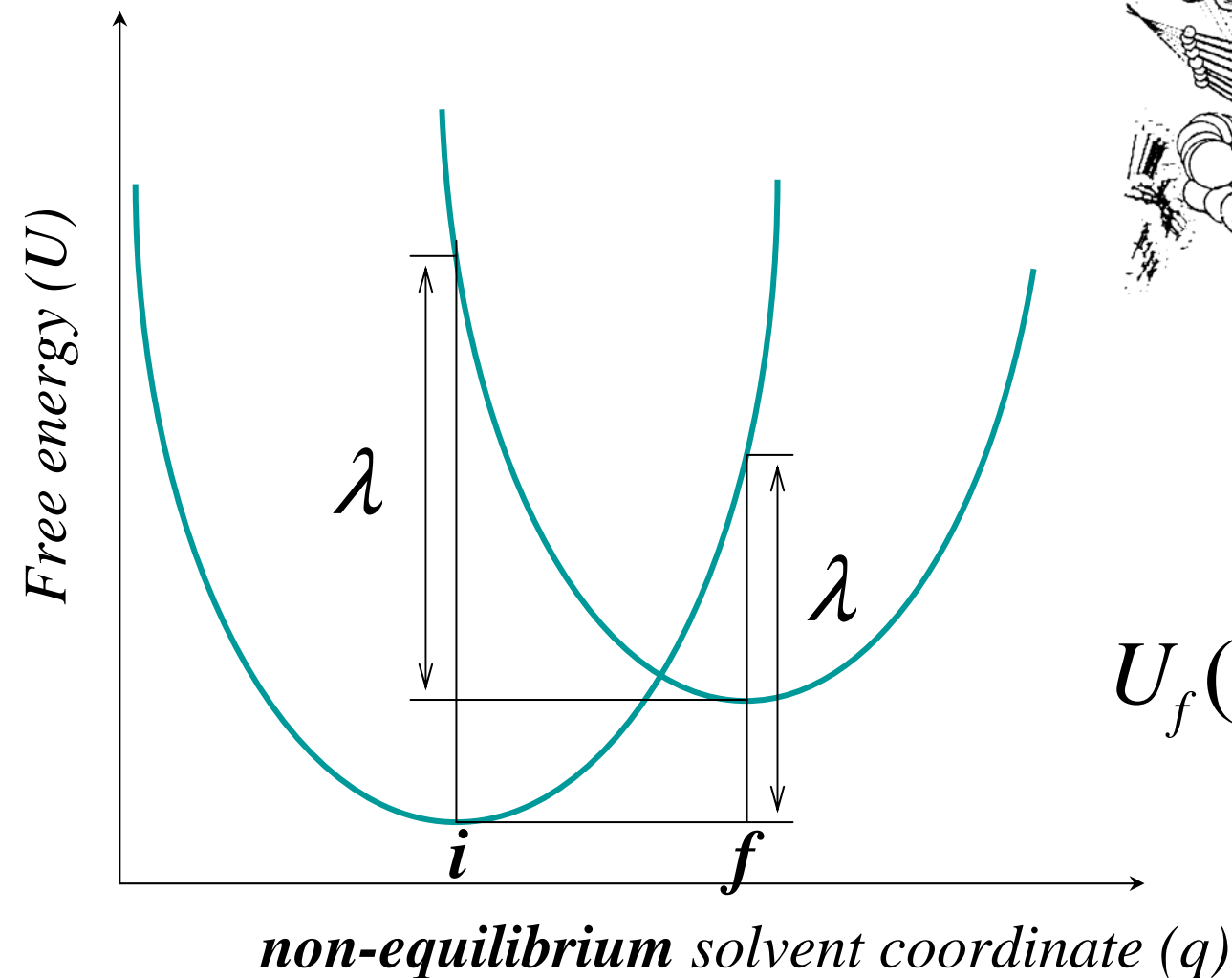
(Ohmine, 1992)

$$U_i(q) = \lambda q^2$$

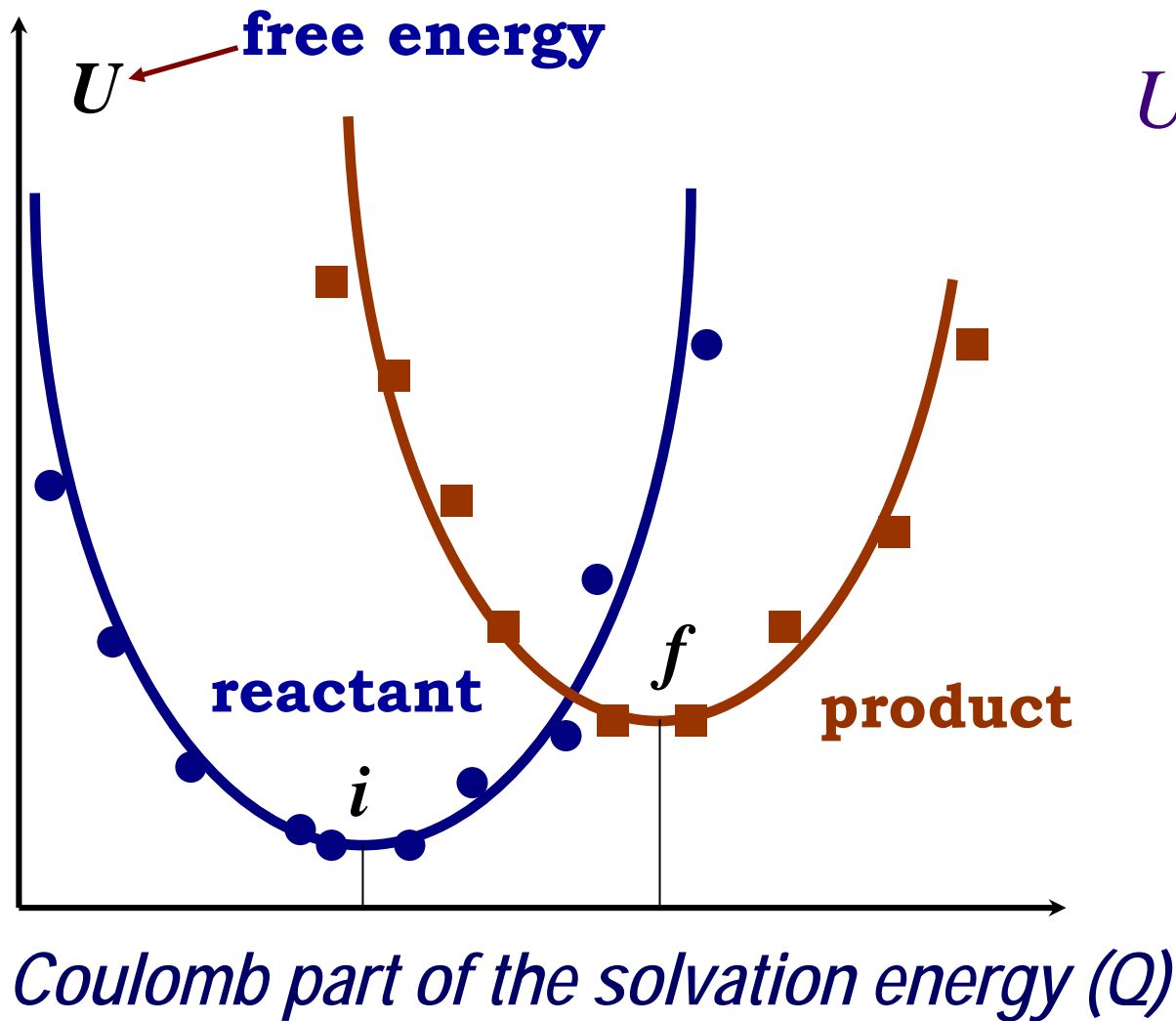
$$U_f(q) = \lambda(q-1)^2 + \Delta F$$

$$\Delta E_a = \frac{(\lambda + \Delta F)^2}{4\lambda}$$

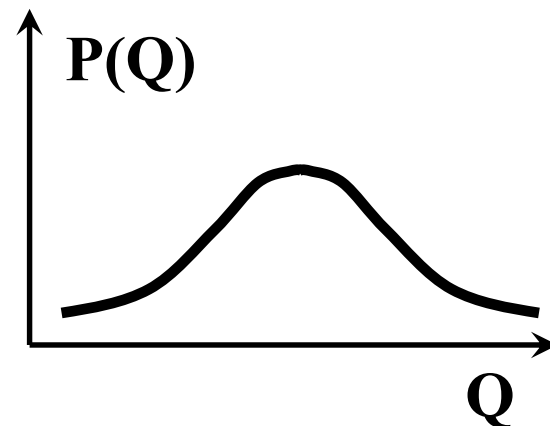
md



Computer simulations



$$U(Q) = -k_B T \ln P(Q)$$



Umbrella sampling

Stochastic theory



Hendrik A. Kramers
/pioneered a stochastic
approach in chemical kinetics/



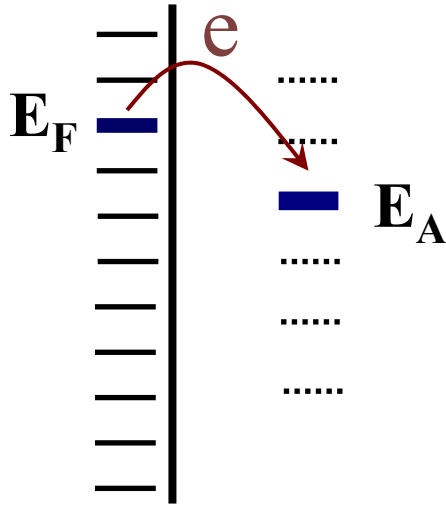
Leonid D. Zusman
/extended Kramers theory
to electron transfer reactions/

Reaction rate depends on dynamical
solvent properties as well
(*friction, viscosity*)



In terms of stochastic theory an overcoming
of the activation barrier more resembles
“climbing” (diffusion)

Solvent correlation function



$$K(\tau) = \langle E_A(0), E_A(\tau) \rangle$$

$$K(\tau) = 2k_B T \lambda_s \frac{Q(\tau)}{Q(0)}$$

$$Q(\tau) = \frac{i}{2\pi} \int_{-\infty}^{\infty} \exp(-i\omega\tau) \left[\frac{1}{\epsilon(\omega)} - \frac{1}{\epsilon_\infty} \right] \frac{d\omega}{\omega}$$

S. Mukamel et al.

dielectric spectrum

spectra

Examples of dielectric spectra

sucrose solutions:

$$\varepsilon(\omega) = \varepsilon_{\infty} + \frac{\Delta\varepsilon_C}{1 + (i\omega\tau_C)^{\alpha}} + \frac{\Delta\varepsilon_D}{1 + i\omega\tau_D}$$

water-EG mixtures:

$$\varepsilon(\omega) = \frac{\Delta\varepsilon_1}{1 + i2\omega\tau_1} + \frac{\Delta\varepsilon_2}{1 + i2\omega\tau_2} + \frac{\Delta\varepsilon_3}{1 + i2\omega\tau_3} + \varepsilon_{\infty}$$

N solvent modes (exact expansion)

Solvent reorganization energy

correlation times

$$K(\tau) = 2k_B T \lambda \sum_{i=1}^N \delta_i \exp(-\tau / \tau_i^*)$$

Solvent correlation function

$$\sum_{i=1}^N \delta_i = 1$$

δ_i is the contribution of i -th mode to the solvent reorganization energy

The solvent reorganization energy is “distributed” among N solvent coordinates.

Reaction free energy surface can be described using N solvent coordinate (q_1, \dots, q_N) and (probably) one intramolecular degree of freedom (r):

$$E_i(q_1, \dots, q_N; r) = \sum_{j=1}^N \delta_j \lambda_j q_j^2 + U_i(r)$$

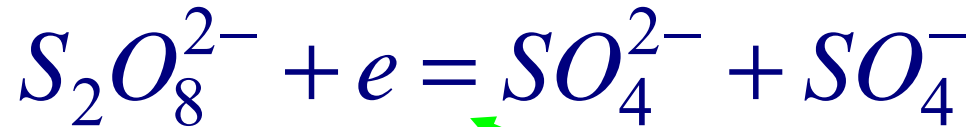
reactant

$$E_f(q_1, \dots, q_N; r) = \sum_{j=1}^N \delta_j \lambda_j (q_j - 1)^2 + U_f(r) + \Delta I$$

product

Usually $N = 2$ (e.g., dimethylacetamide), 3 (EG, alcohols etc)

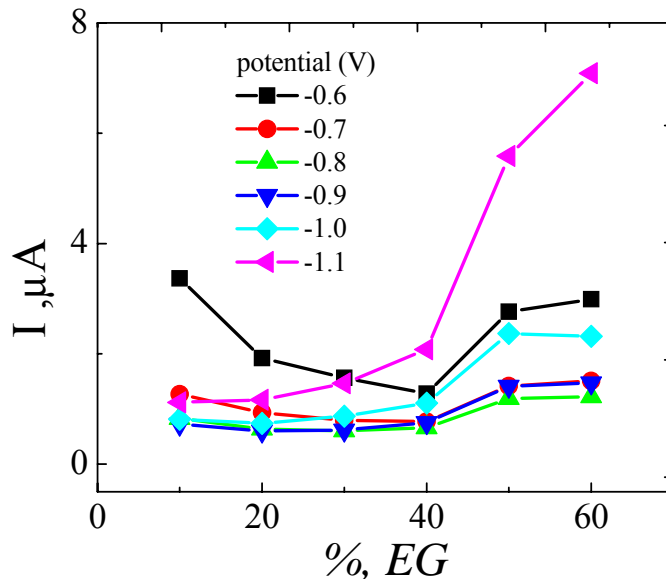
$S_2O_8^{2-}$ reduction at a mercury electrode from water-EG mixtures



- reaction is adiabatic

The first ET is rate limiting

- BBET reaction proceeds at large overvoltages, in the vicinity of activationless discharge, i.e. at small activation barriers

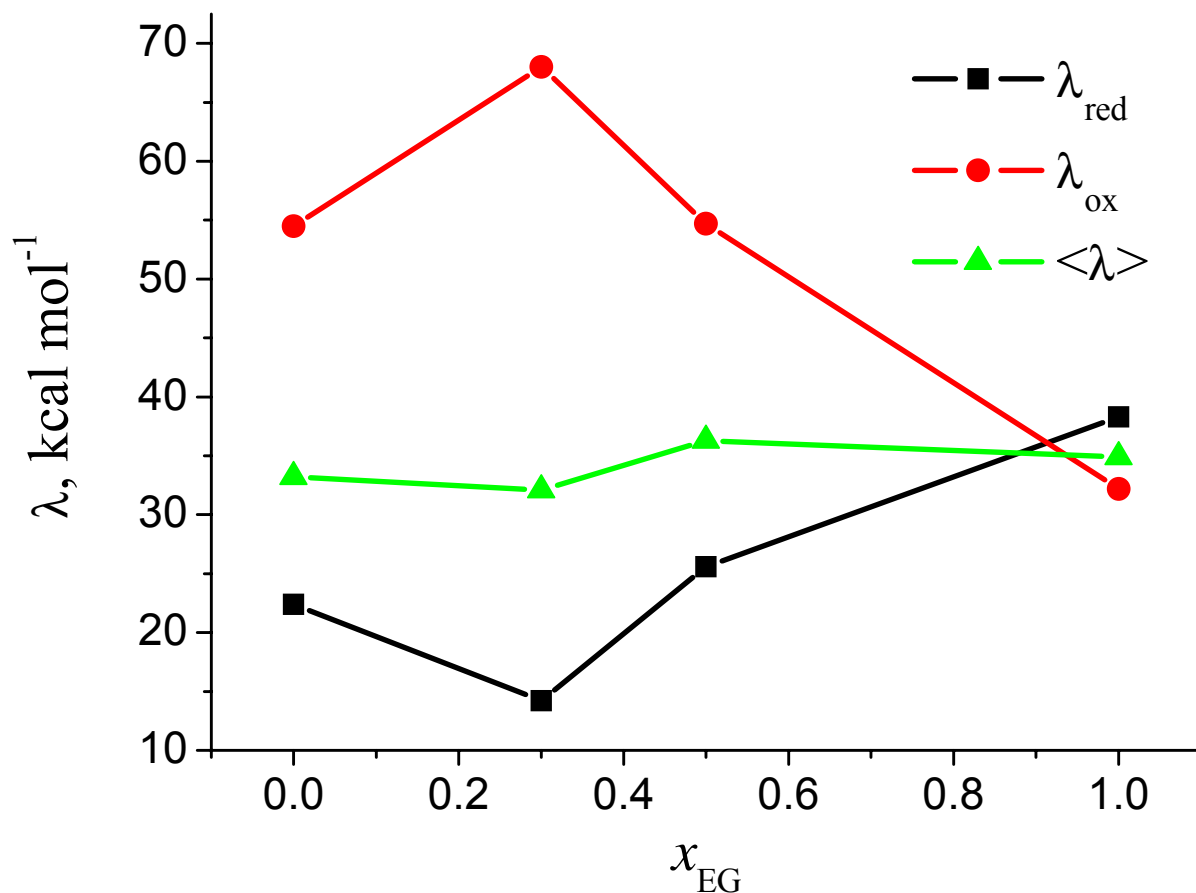


- reaction reveals an anomalous solvent viscosity effect

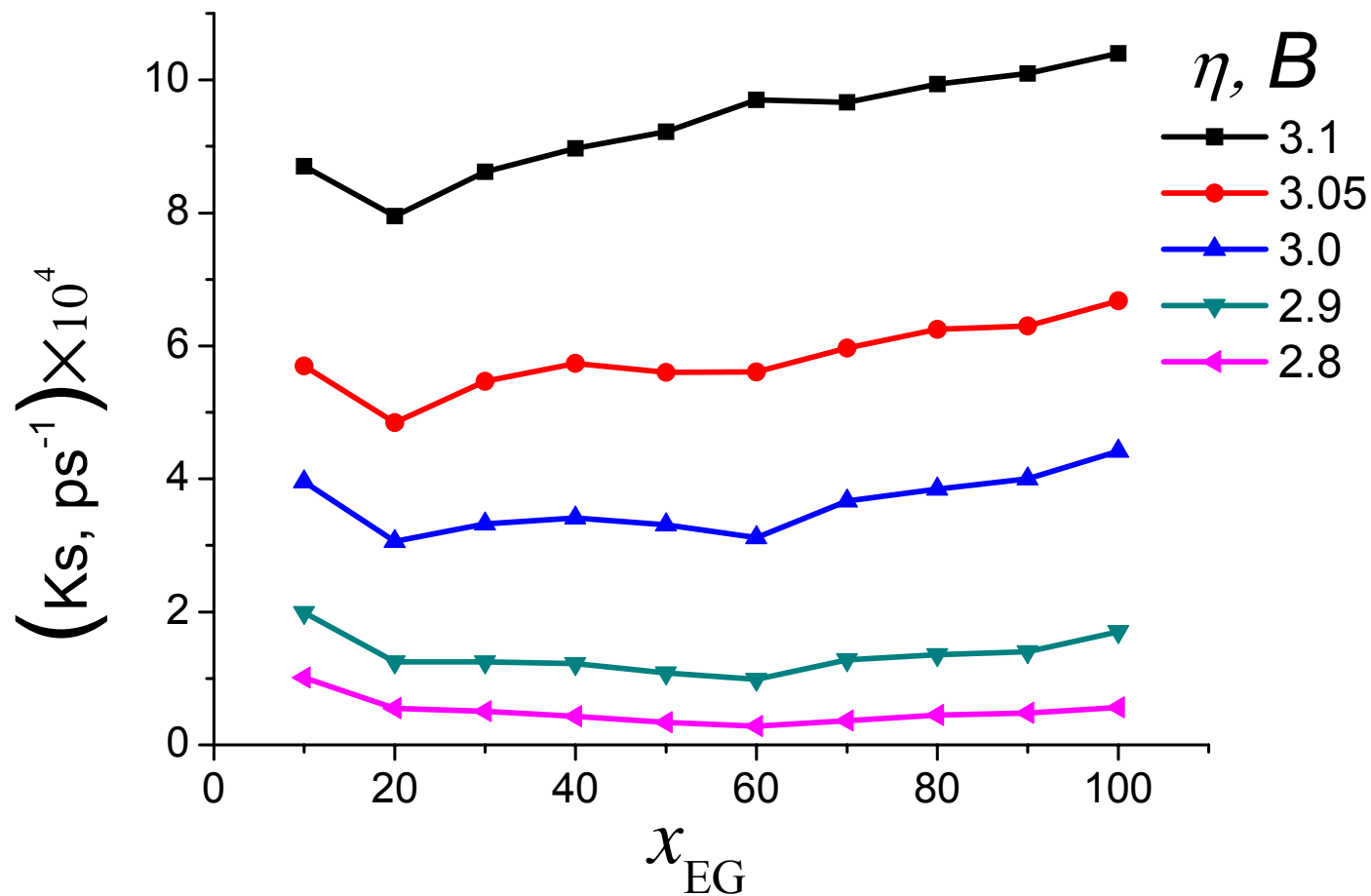
Exp. data (P.A. Zagrebina et al)

- Pekar factor in the solvent reorganization energy is nearly constant
- MD simulations predict even a slight increase of $\langle \lambda \rangle$

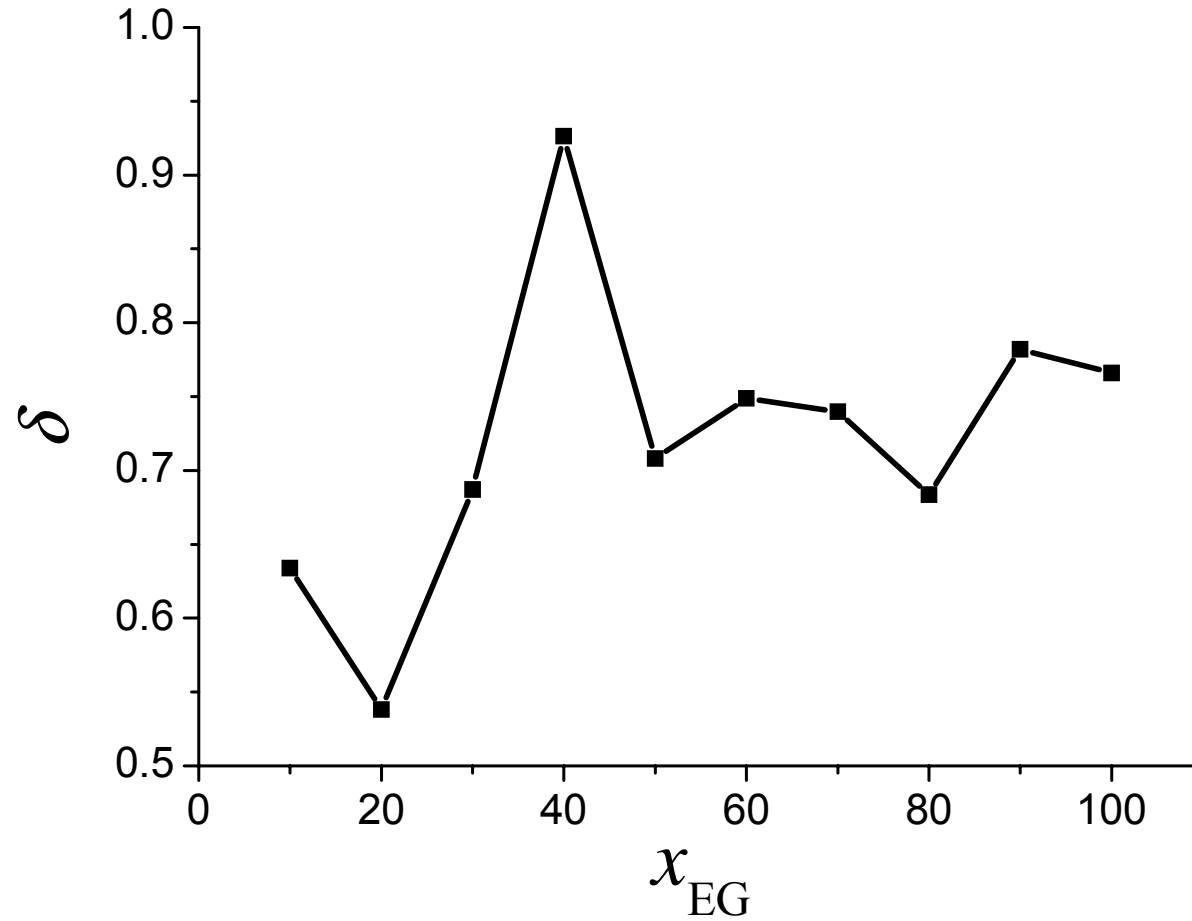
Bulk contributions to the solvent reorganization energy as computed from molecular dynamics (O. Ismailova, M. Probst et al)



Results of Langevin molecular dynamics simulations



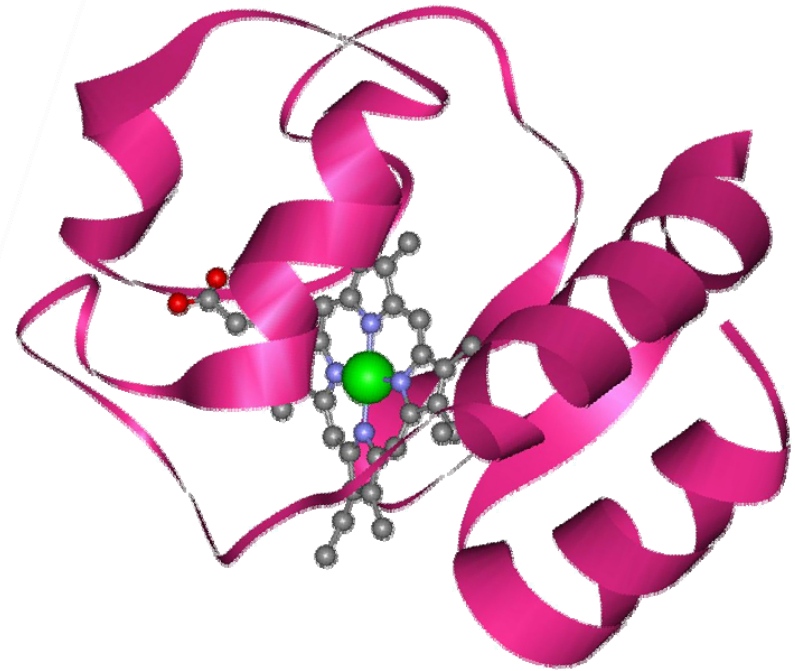
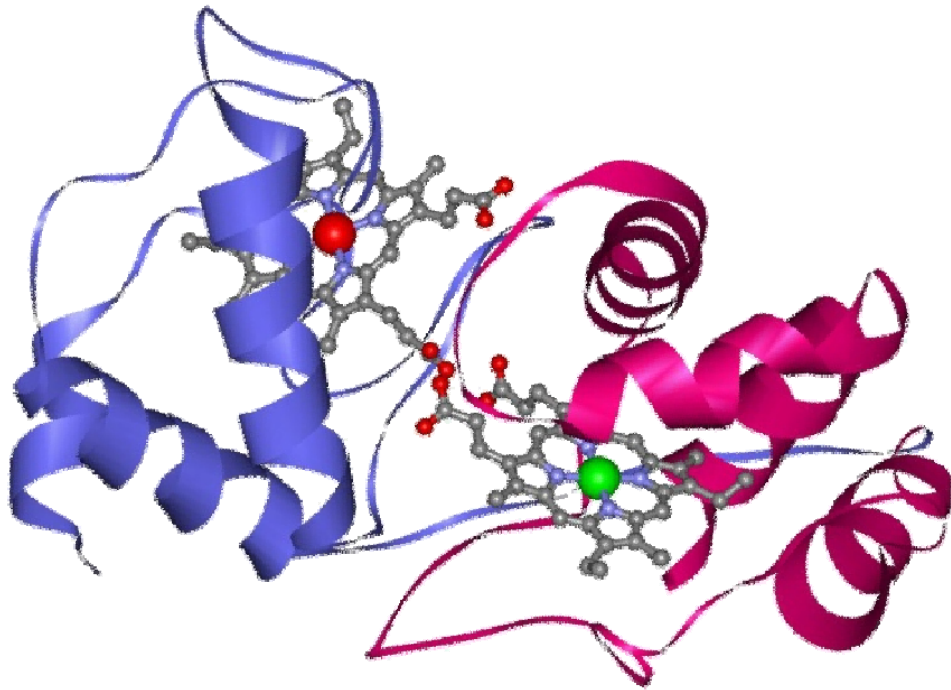
An attempt to explain: saddle point avoidance



Non-Gaussian fluctuations

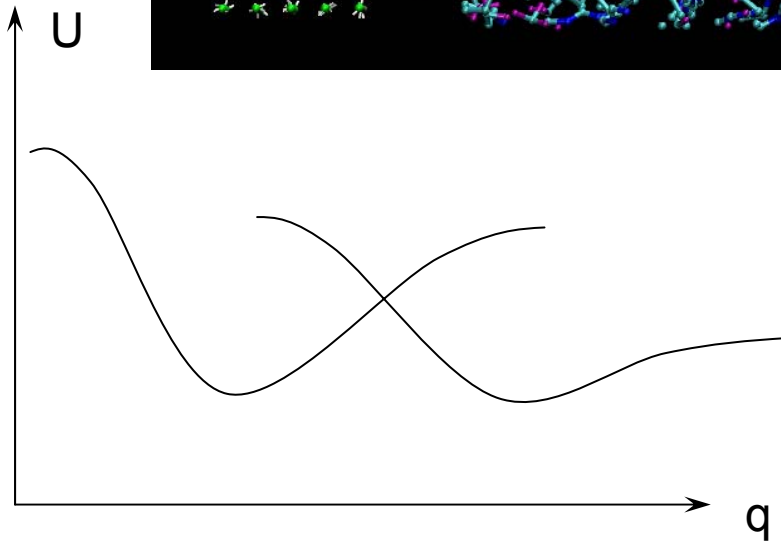
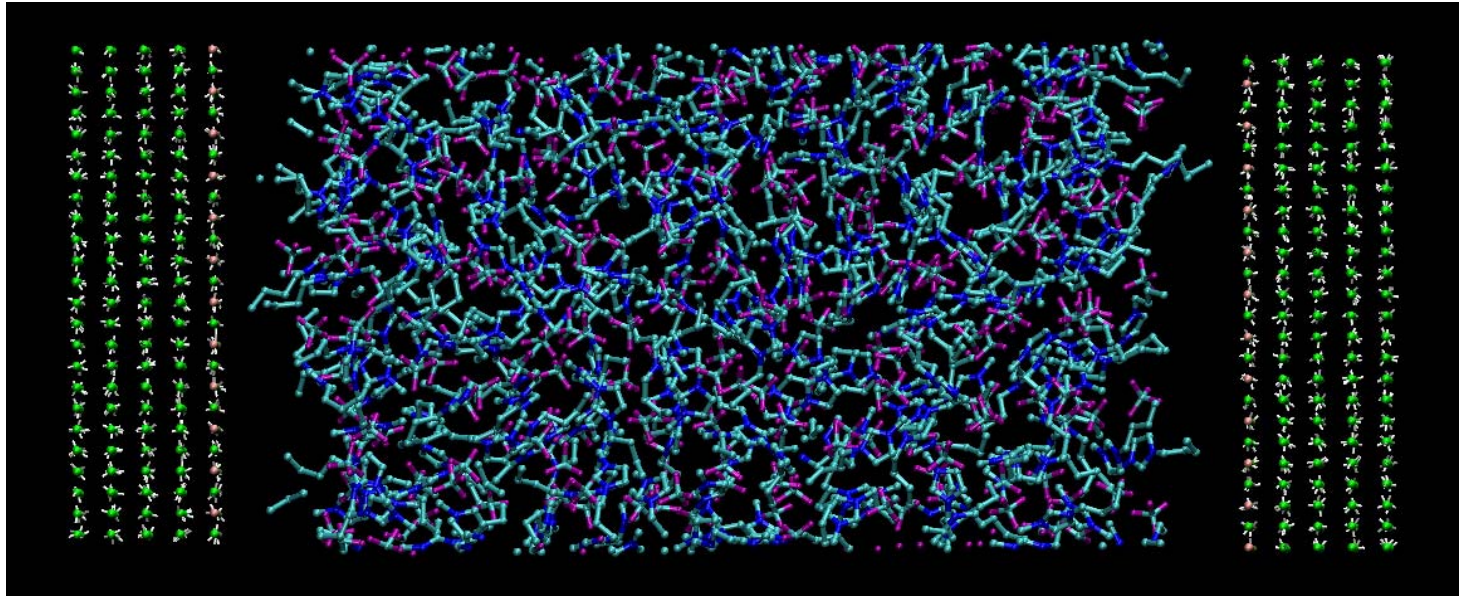
- ferroelectric domains at a protein/water interface

D.N. LeBard, D.V. Matyushov, PCCP, 12 (2010) 15335



MD simulation of the Au(111)/[BMIM][BF₄] interface

(S.A. Kislenko et al)

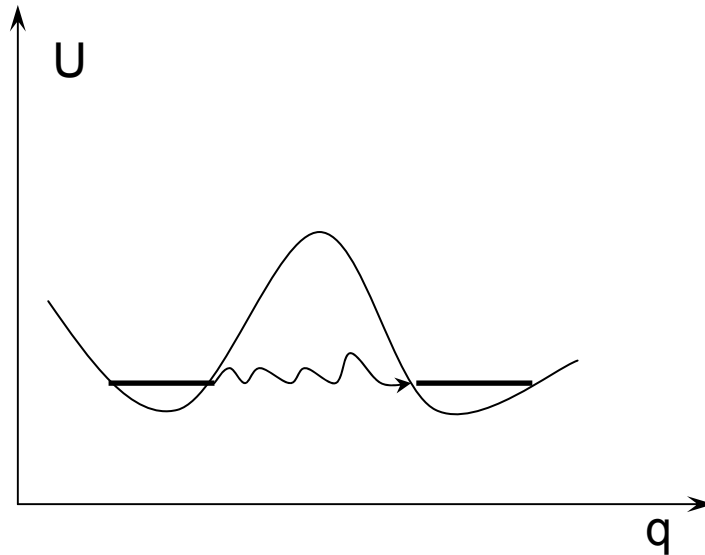
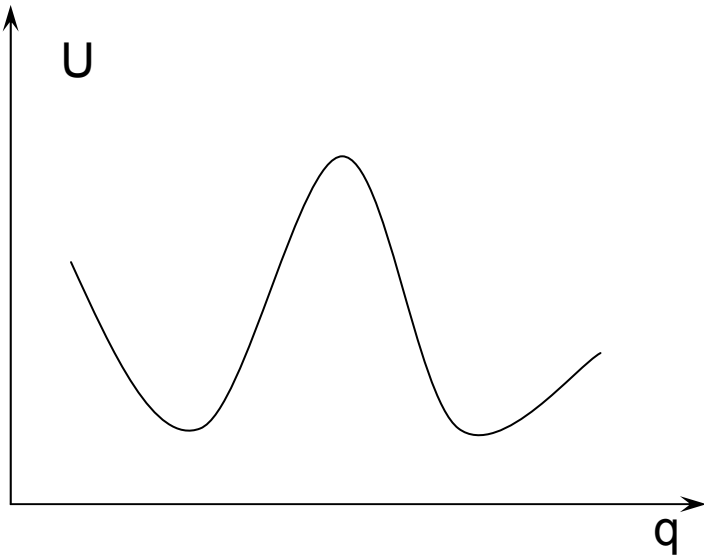


Non-linear response ?

quant

Solvent coordinate vs Quantum effects

- decreasing of the activation barrier → increasing rate constant
- tunneling → decreasing rate constant



Effect of solvent quantum modes

$$k = \exp\left[-\frac{\Delta E_a^*}{k_B T}\right] \exp[-\sigma] = \exp\left[-\frac{(\lambda_s^* - \eta)^2}{4\lambda_s^* k_B T}\right] \exp[-\sigma]$$

$$\lambda_s^* = \xi \lambda_s$$

$$\xi = \frac{2}{\pi C} \int_0^{\omega^*} \frac{\text{Im } \varepsilon(\omega)}{\omega \|\varepsilon(\omega)\|^2} d\omega$$

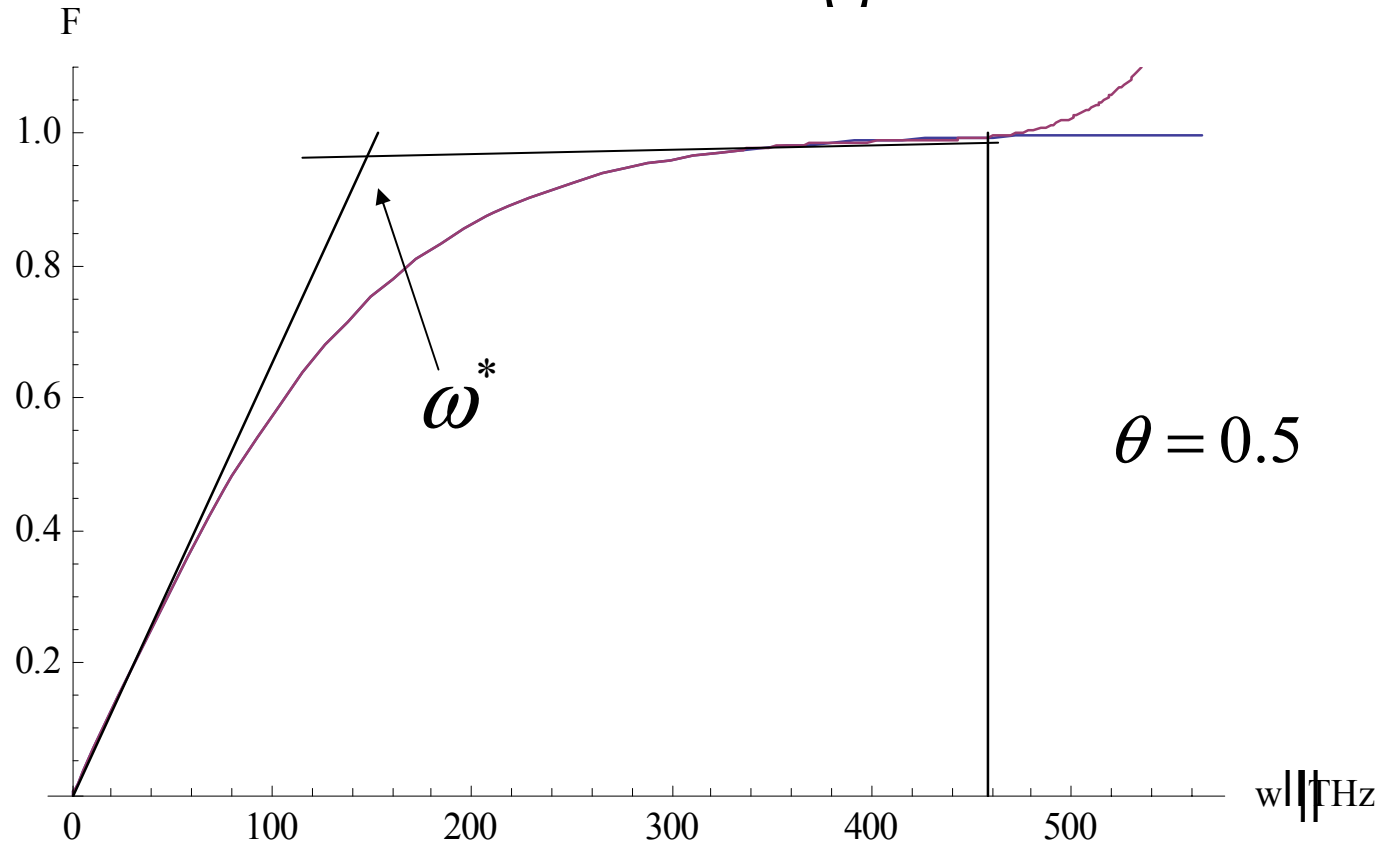
$$C = \frac{1}{\varepsilon_\infty} - \frac{1}{\varepsilon_{st}}$$

$$\sigma = \frac{2\lambda_s}{\pi C} \int_{\omega^*}^{\infty} \frac{\text{Im } \varepsilon(\omega)}{\omega^2 \|\varepsilon(\omega)\|^2} d\omega$$

tunneling factor

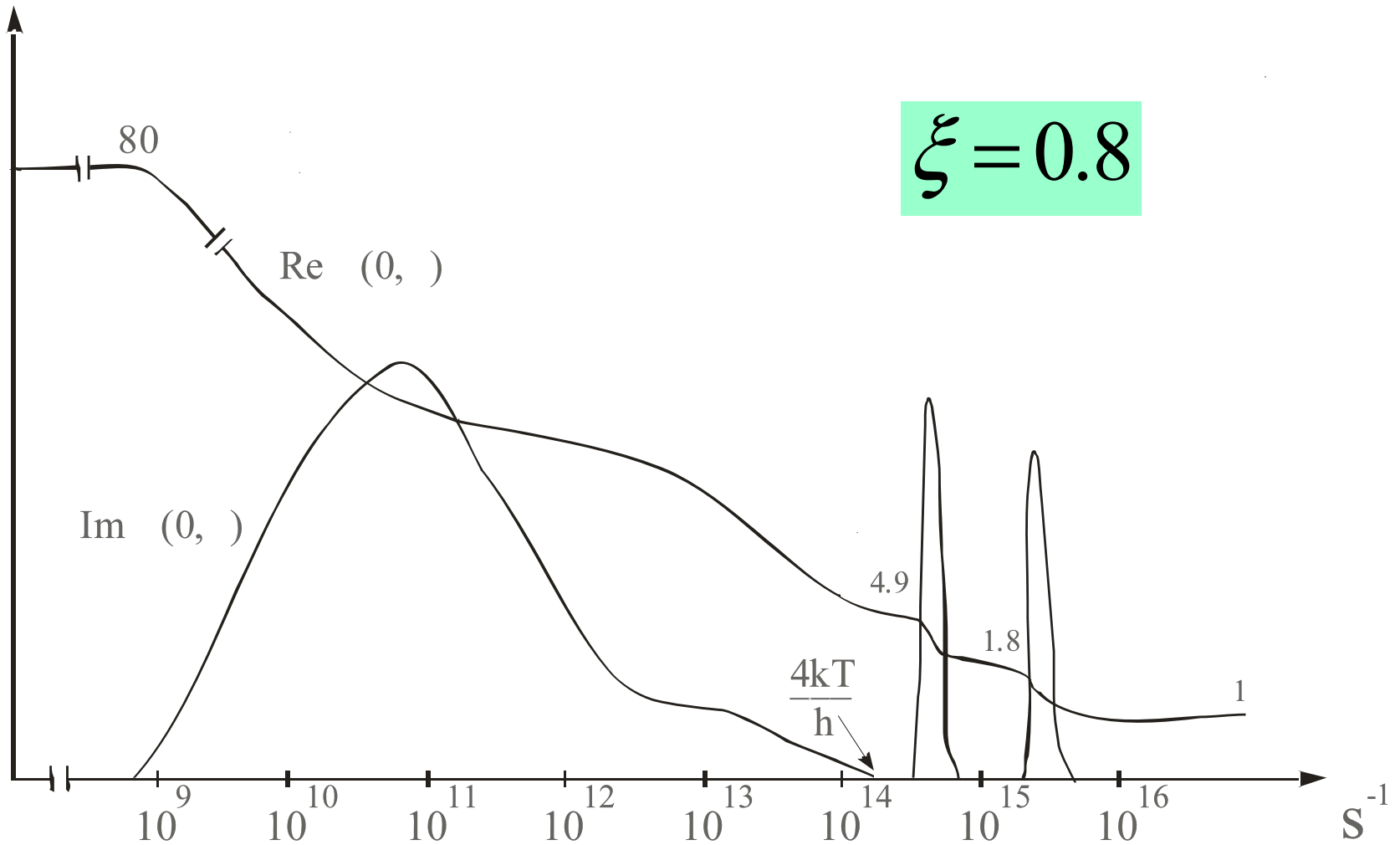
Pekar factor

Interpolation ch...sh w, w^* by polynomial of degree 6
 on the interval $0, 3w^*$

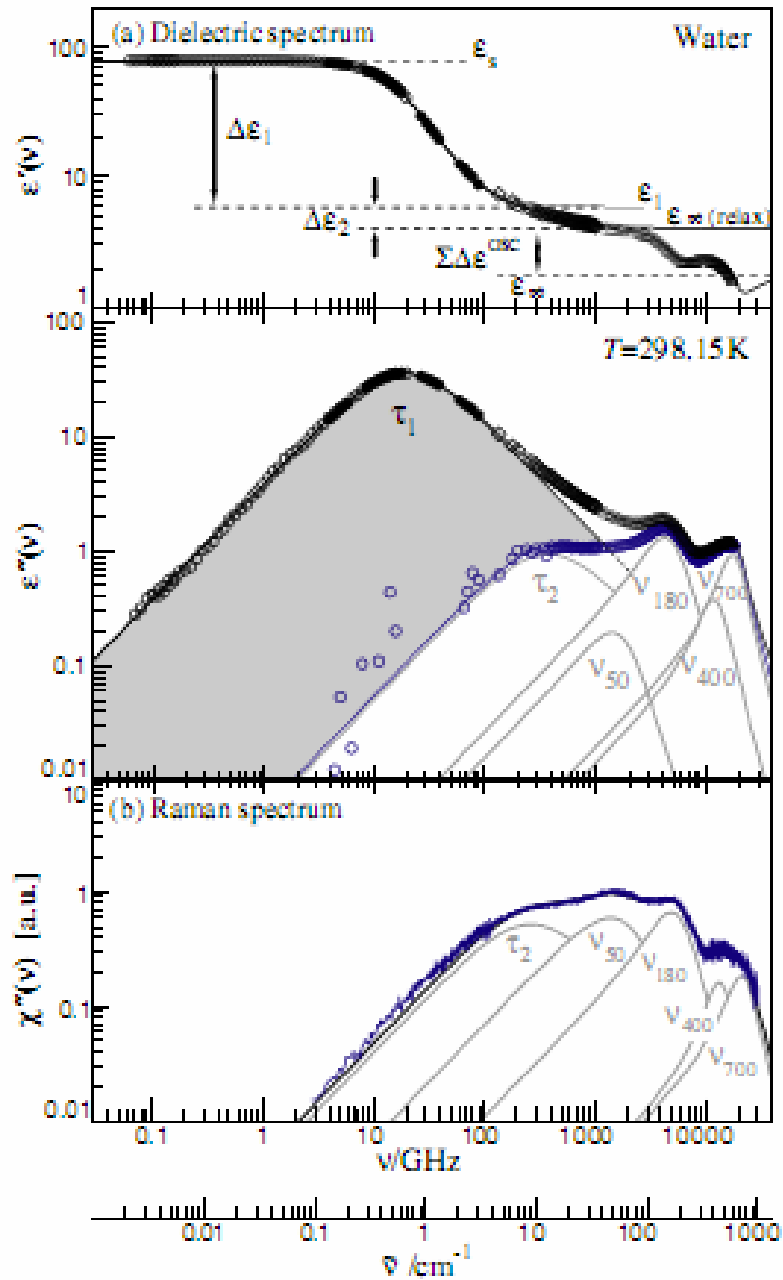


$$F(\theta, \omega) = \frac{ch\{\beta \ \omega / 2\} - ch\{(1 - 2\theta)\beta \ \omega / 2\}}{sh\{\beta \ \omega / 2\}}$$

Dielectric spectra of water (J.A. Saxton, 1953)



R. Buchner and
co-workers (2005)

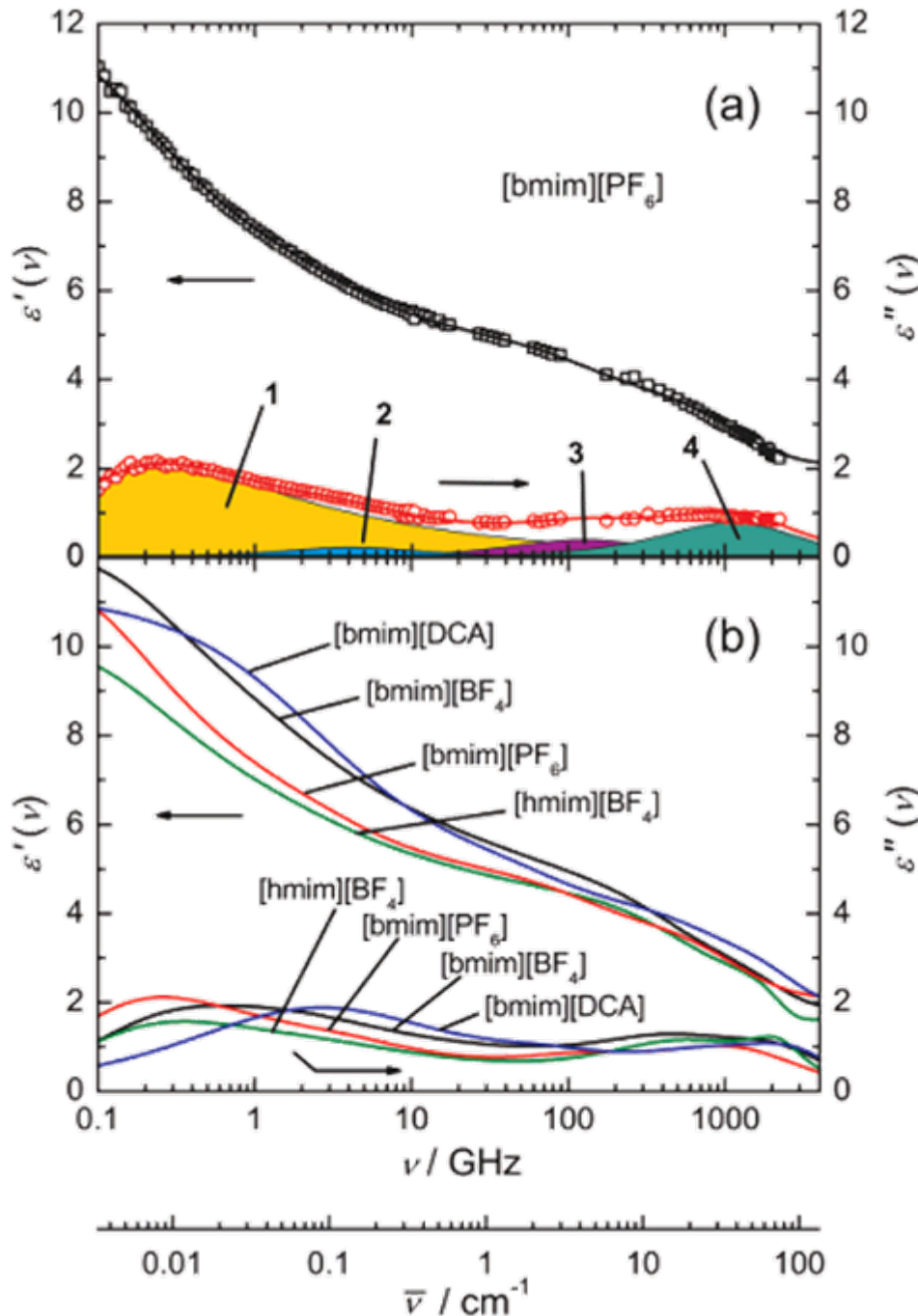


$$\begin{aligned}
 \epsilon(\nu) = & \epsilon_{\infty} + \frac{S_1}{1+i2\pi\nu\tau_1} + \frac{S_2}{1+i2\pi\nu\tau_2} + \\
 & \frac{S_3}{1+i2\pi\nu\tau_3} + \frac{S_4\nu^2}{\nu_4^2 - \nu^2 + i\gamma_4\nu} \\
 & + \frac{S_5\nu^2}{\nu_5^2 - \nu^2 + i\gamma_5\nu} + \frac{S_6\nu^2}{\nu_6^2 - \nu^2 + i\gamma_6\nu} + \\
 & + \frac{S_7\nu^2}{\nu_7^2 - \nu^2 + i\gamma_7\nu}
 \end{aligned}$$

$\xi = 0.9$

R. Buchner and
co-workers (2008)

Dielectric spectra of some ionic liquids



$$\xi = 0.55$$

Electronic transmission coefficient

$$\kappa_e = \frac{1 - \exp(-2\pi\gamma_e)}{1 - (1/2)\exp(-2\pi\gamma_e)}$$

Landau-Zener factor

half of resonance splitting

$$\gamma_e = \left(\frac{\Delta E_e}{2} \right)^2 \sqrt{\frac{\pi}{(\lambda_s + \lambda_{in}) \cdot k_B T}}$$

effective frequency

Two important limiting cases:

$$\gamma_e \ll 1 \Rightarrow \kappa_e \approx \gamma_e \quad \text{(non-adiabatic)}$$

$$\gamma_e \gg 1 \Rightarrow \kappa_e \approx 1 \quad \text{(adiabatic)}$$

It is reasonable to employ the perturbation theory for large molecular systems

$$\frac{\Delta E_e}{2} \approx \int \Psi_i \hat{V} \Psi_f dV - \int \Psi_i \hat{V} \Psi_i dV \cdot \int \Psi_i \Psi_f dV$$

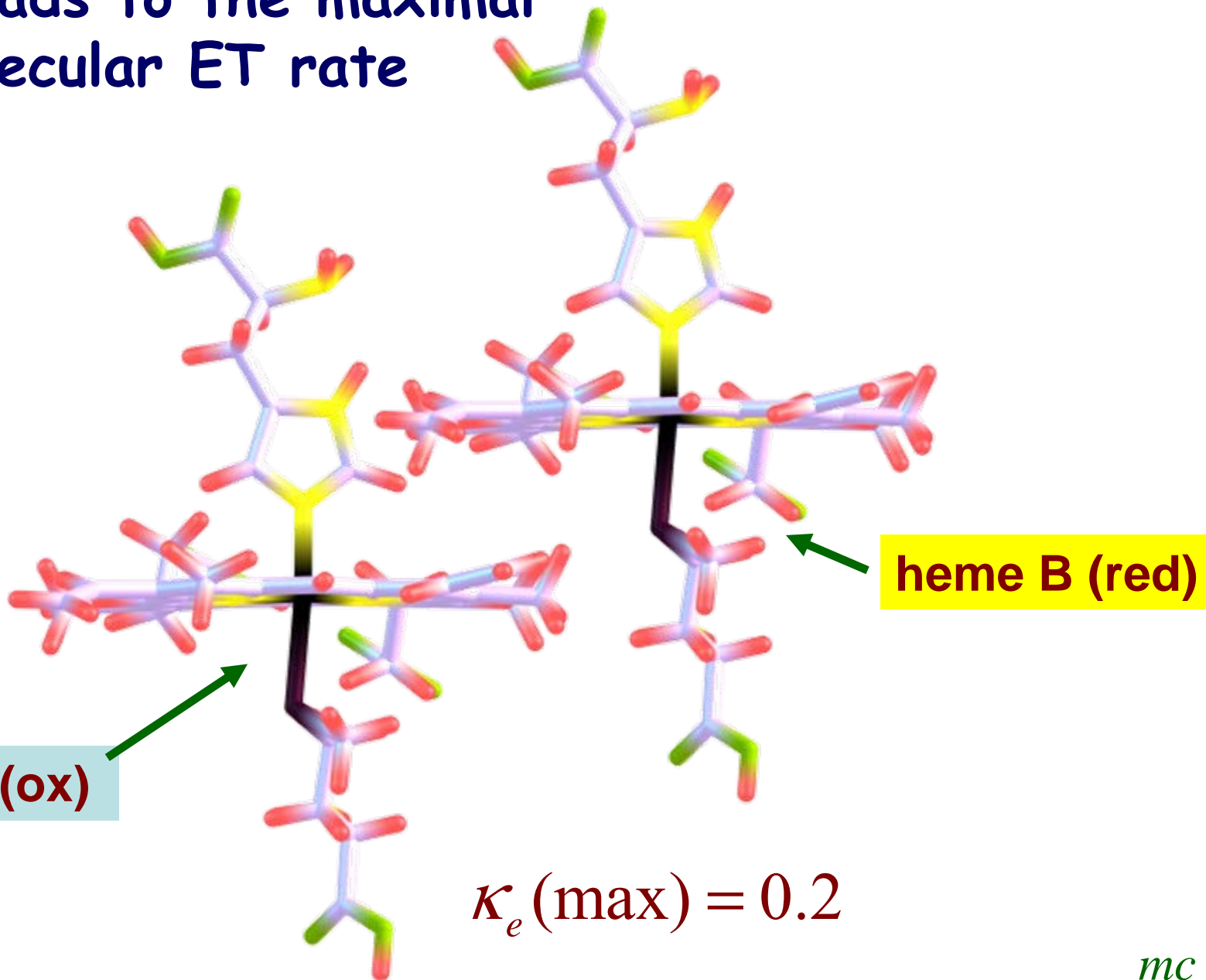
Perturbation (molecular electrostatic potential)

$$V(r) = -\sum_i \frac{Z_i}{|R_i - r|} + \sum_j \int \frac{|\psi_j(r')|^2 d\Omega'}{|r' - r|}$$

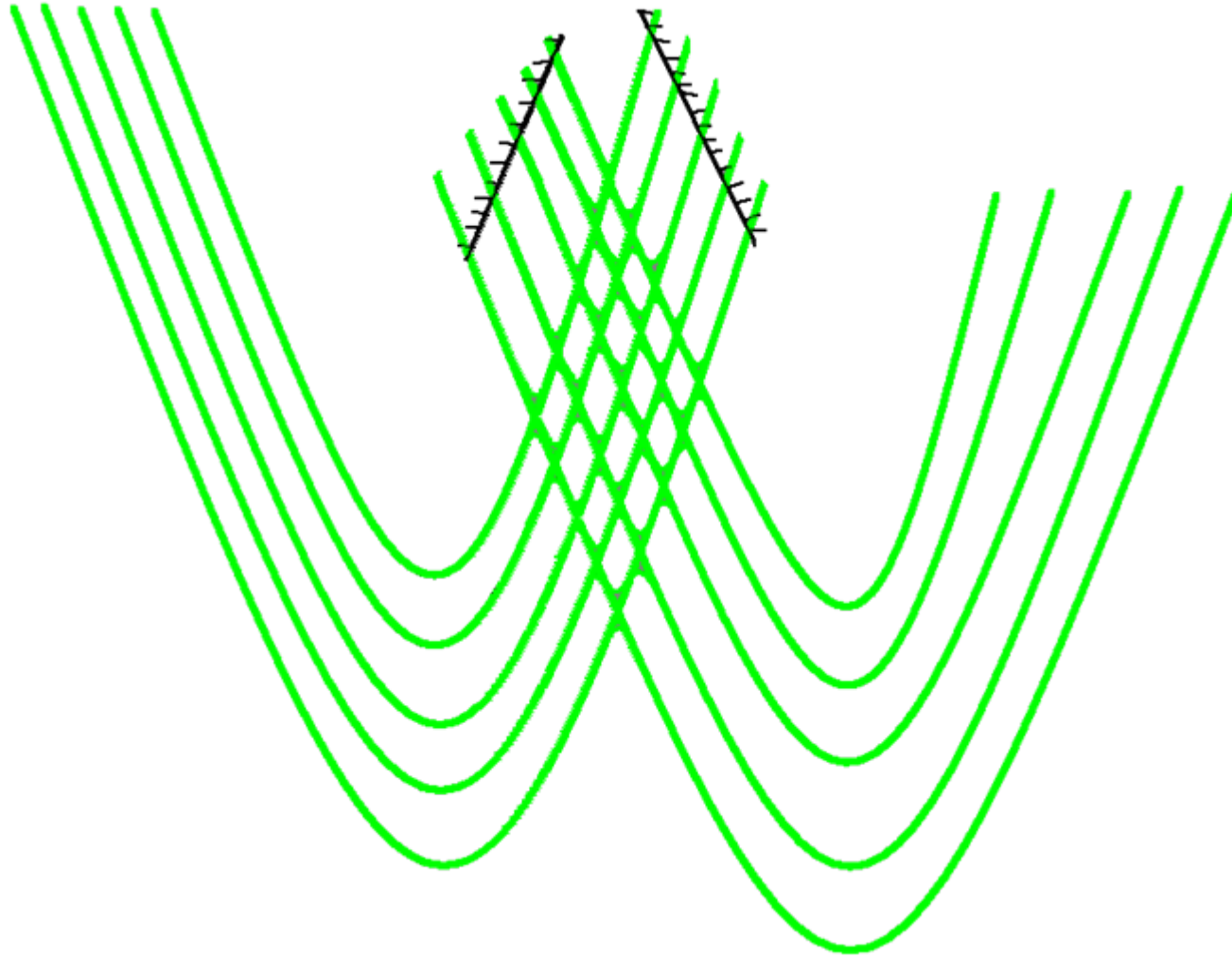
$$V(r) \approx \sum_i \frac{q_i^*}{|R_i - r|}$$

ChelpG atomic charges

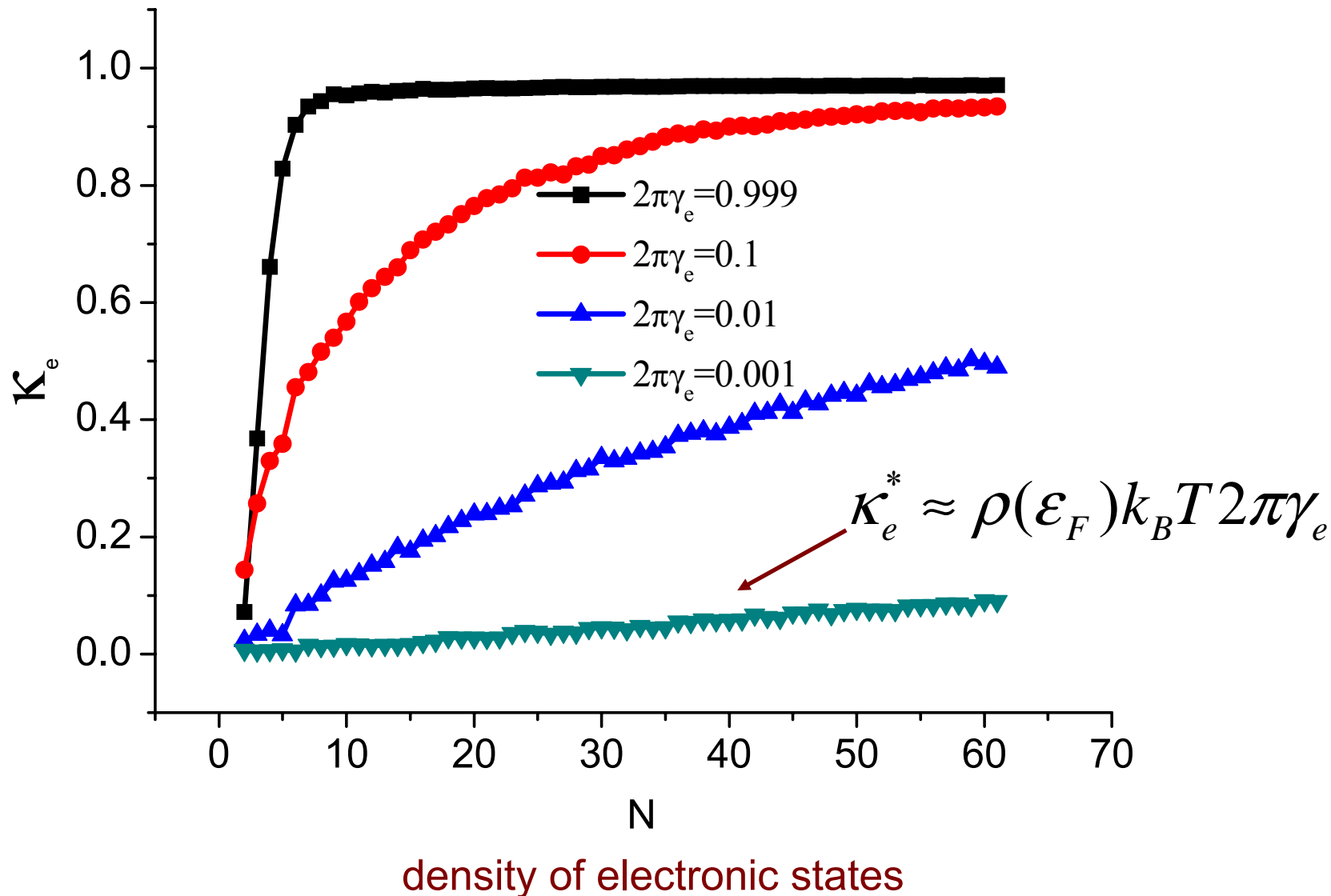
Orientation of the cyt c_4 heme groups which leads to the maximal intramolecular ET rate



Моделирование методом Монте-Карло (случайное блуждание по узлам двумерной решётки)

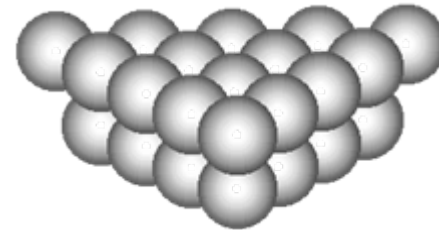
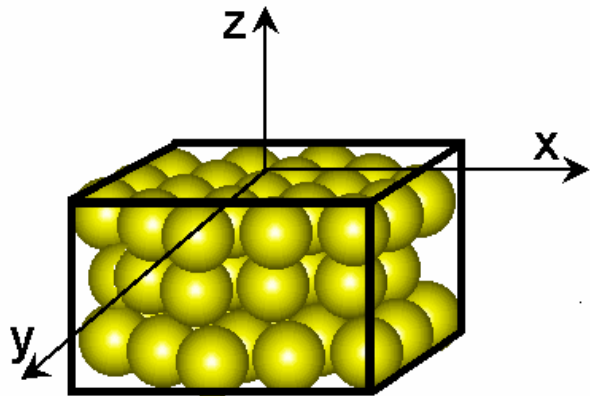


Electronic transmission coefficient vs density of electronic states calculated with the help of MC simulations at different values of γ_e



Model calculation of electronic transmission coefficient for interfacial reactions: some challenges.

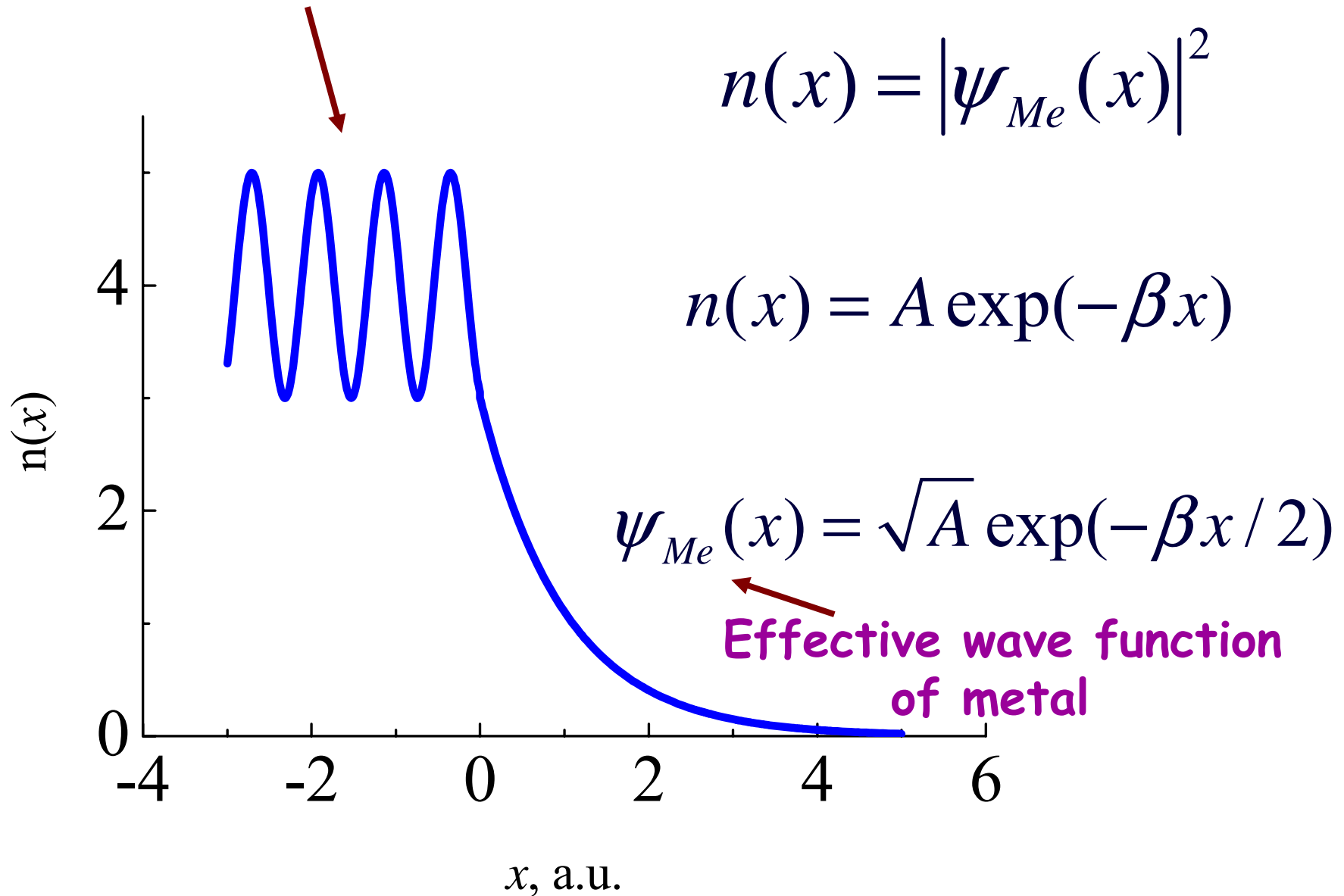
1. Model of a charged metal surface
(cluster, slabs, “jellium”, etc)



2. Solvent effect on the wave functions and perturbation

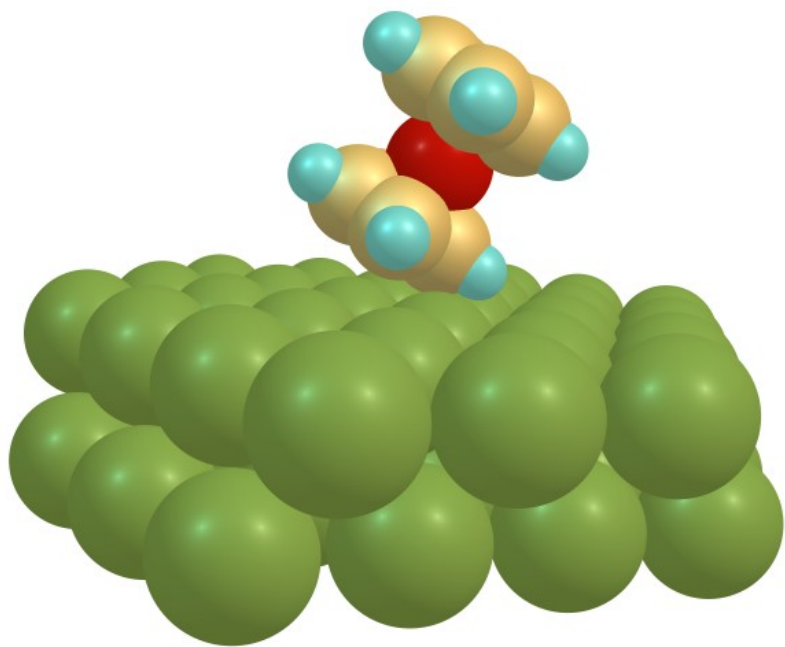
3. Asymptotic behaviour of wave functions

Electronic density of metal slabs vs distance

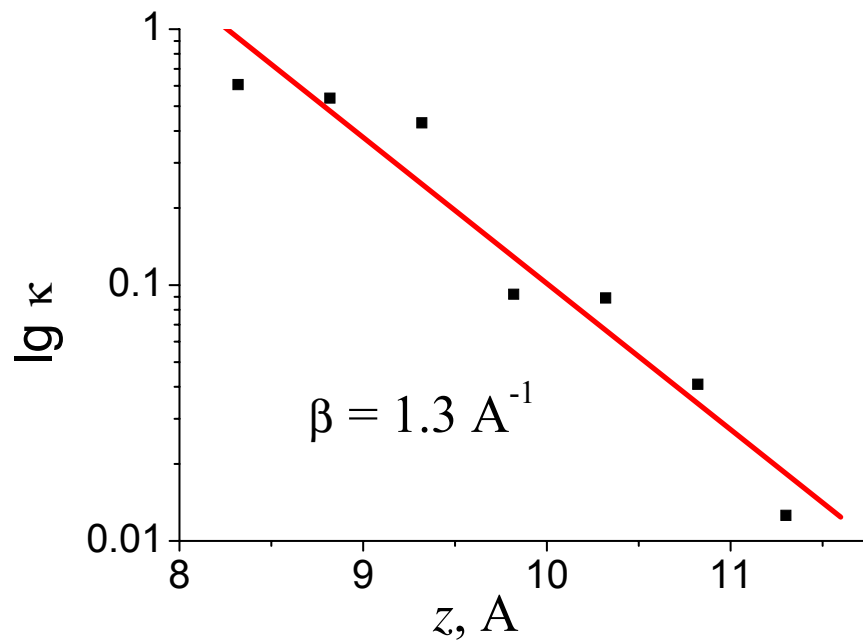




$$\kappa = \kappa_0 \exp(-\beta z)$$

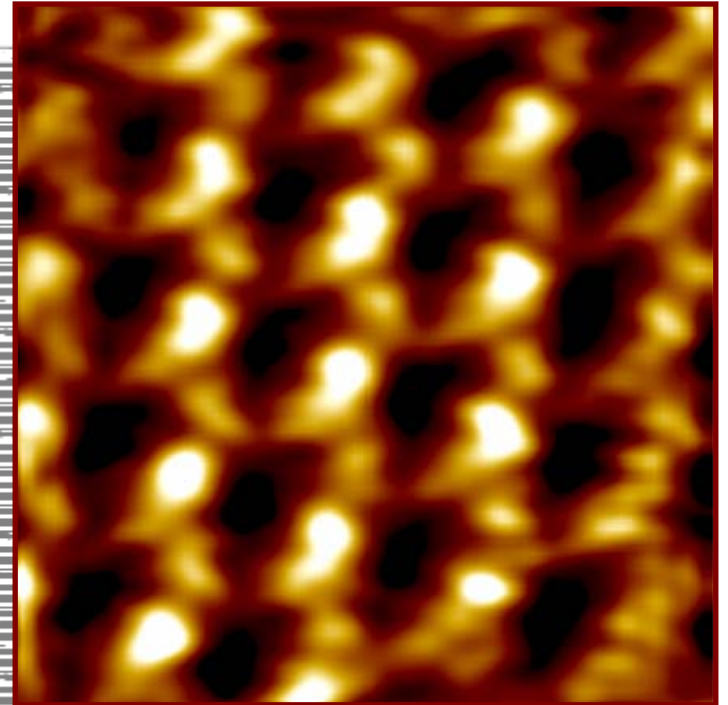
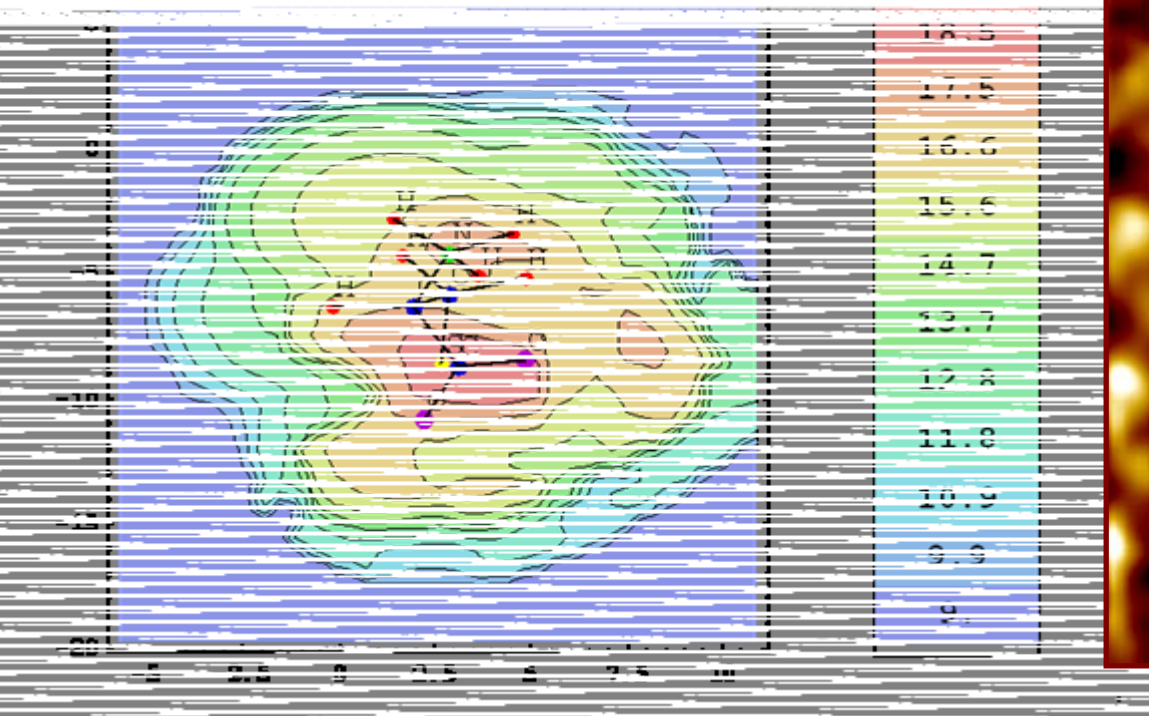


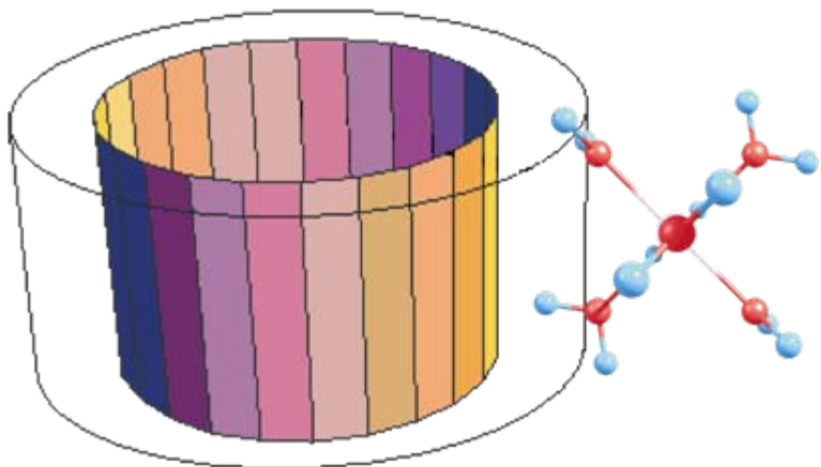
Au₅₄ cluster



Model STM contrast

Cysteine adsorption on
Au(110) electrode
(in situ STM images)





Me(111) vs monoatomic wires Fe(III)/Fe(II)

